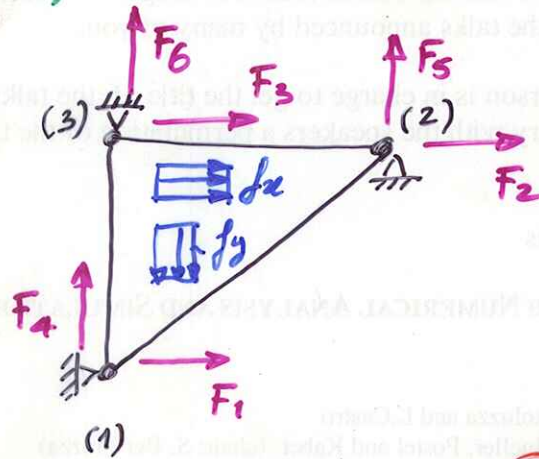


VECTOR DAS FORÇAS NODAIS EQUIVALENTES

• FORÇAS DE MASSA •



$$\underline{F} = \int_V \underline{\psi}^T \underline{f} dv = \int_V \begin{bmatrix} \psi_1 & 0 \\ \psi_2 & 0 \\ \psi_3 & 0 \\ 0 & \psi_1 \\ 0 & \psi_2 \\ 0 & \psi_3 \end{bmatrix} \begin{bmatrix} f_x \\ f_y \end{bmatrix} dv$$

$$\underline{F} = \int_V \begin{bmatrix} \psi_1(x,y) f_x(x,y) \\ \psi_2(x,y) f_x(x,y) \\ \psi_3(x,y) f_x(x,y) \\ \psi_1(x,y) f_y(x,y) \\ \psi_2(x,y) f_y(x,y) \\ \psi_3(x,y) f_y(x,y) \end{bmatrix} dv$$

CASO PARTICULAR



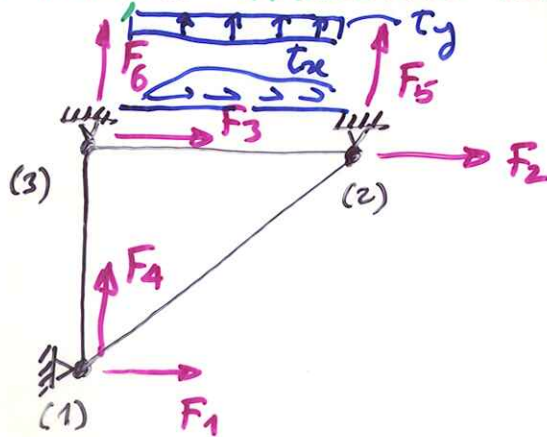
+ FORÇAS DE MASSA CONSTANTES

$$\underline{F} = \frac{1}{3} \begin{bmatrix} A f_x \\ A f_x \\ A f_x \\ A f_y \\ A f_y \\ A f_y \end{bmatrix}$$

A = área do elemento

VECTORES DAS FORÇAS NODAIS EQUIVALENTES

FORÇAS APLICADAS NA FRONTEIRA.

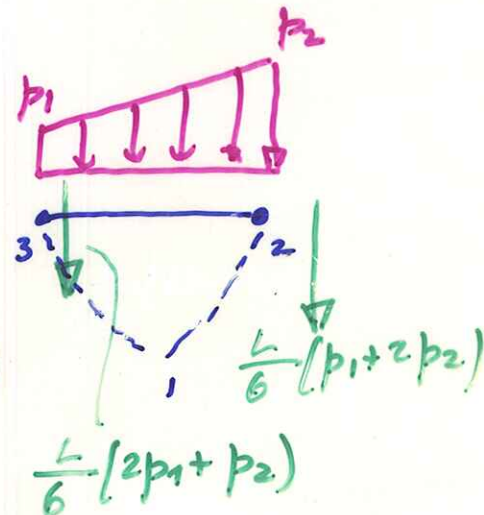
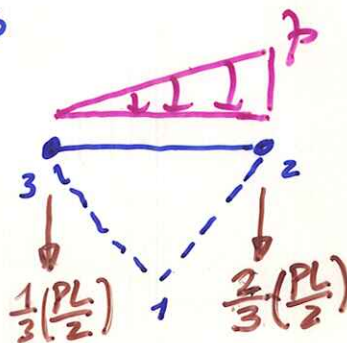
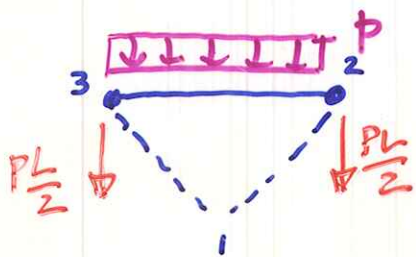


$$\underline{F} = \int_r \bar{\Psi} \underline{T} dr = \int_r \begin{bmatrix} \bar{\Psi}_1 & 0 \\ \bar{\Psi}_2 & 0 \\ \bar{\Psi}_3 & 0 \\ 0 & \bar{\Psi}_1 \\ 0 & \bar{\Psi}_2 \\ 0 & \bar{\Psi}_3 \end{bmatrix} \begin{bmatrix} t_x \\ t_y \end{bmatrix} dr$$

INTEGRAÇÃO AO LONGO DO LADO CARREGADO

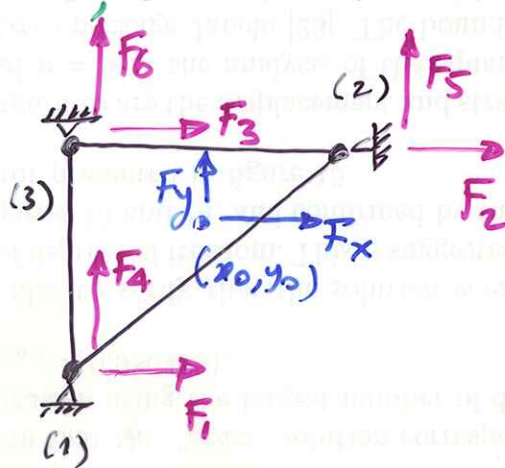
$$\underline{F} = \int_r \begin{bmatrix} \bar{\Psi}_1(x,y) t_x \\ \bar{\Psi}_2(x,y) t_x \\ \bar{\Psi}_3(x,y) t_x \\ \bar{\Psi}_1(x,y) t_y \\ \bar{\Psi}_2(x,y) t_y \\ \bar{\Psi}_3(x,y) t_y \end{bmatrix} dr$$

CASOS PARTICULARES



VECTORES DAS FORÇAS NODAIS EQUIVALENTES

• FORÇAS CONCENTRADAS •



$$\underline{F} = \begin{bmatrix} F_x \psi_1(x_0, y_0) \\ F_x \psi_2(x_0, y_0) \\ F_x \psi_3(x_0, y_0) \\ \hline F_y \psi_1(x_0, y_0) \\ F_y \psi_2(x_0, y_0) \\ F_y \psi_3(x_0, y_0) \end{bmatrix}$$